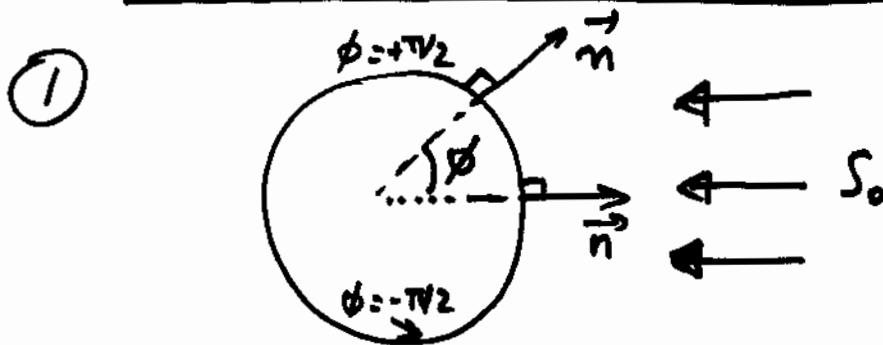




at a latitude ϕ and longitude λ
the projection of the solar energy flux
onto the local normal ($\vec{n}$) is:

$$S_0 \cos \phi \cos \lambda$$



The infinitesimal area element at (λ, ϕ) is
 $dS = (R \cos \phi d\lambda) \times R d\phi$

$\therefore$ energy received by Earth per
unit time is



$$\begin{aligned} & \int_{-\pi/2}^{+\pi/2} S_0 \cos \lambda \int_{-\pi/2}^{+\pi/2} \cos \phi (R^2 \cos \phi d\phi d\lambda) \\ &= R^2 S_0 \int_{-\pi/2}^{+\pi/2} \cos \lambda d\lambda \int_{-\pi/2}^{+\pi/2} \cos^2 \phi d\phi = \pi R^2 S_0 \\ &\sim 175 \text{ PW} = 175,000 \text{ TW!} \end{aligned}$$

② $T_E \equiv \left[\frac{(1-\alpha_p) S_0}{4\sigma} \right]^{1/4}$ standard values:
 $S_0 = 1,368 \text{ Wm}^{-2}$ $\alpha_p = 0.3$

(i) with S_0 ↓ by 30% but keeping $\alpha_p = 0.3$

$$T_E \sim \left[\frac{(1-0.3) 910}{4 \times 5.7 \cdot 10^{-8}} \right]^{1/4} \sim 230 \text{ K}$$

The Earth was much warmer back then because, one believes, the atmospheric CO_2 concentration was much larger than present. The surface gained more energy through atmospheric

radiation from above than from the sun.

- (ii) $\alpha_p = 0.1 \therefore T_E \sim 271\text{K}$. This is below freezing so the "Tern model" can not be used for this calculation (it leads to an inconsistency).

③

Top of Atmosphere: $\frac{(1-\alpha_p)S_0}{4} = \sigma T_A^4$

surface: $\sigma T_A^4 = \sigma T_s^4$

(downward arrow) = (upward arrow)

$\therefore T_A = T_s = T_E$! Even with a perfect absorbing layer the surface is the same as that obtained with the "Tern model"! What matters for the greenhouse effect is to absorb radiation from the surface AND leave un-attenuated the solar radiation.

④ (i) $\Gamma = \frac{303-255}{5} = 9.6\text{ K/km}$ (ii) $\Gamma = \frac{288-255}{5} = 6.6\text{ K/km}$

The first temperature profile is most likely to be unstable.

⑤

Top of Atm. $\left\{ \begin{array}{l} \frac{(1-\alpha_p)S_0}{4} = \sigma T_A^4 \Rightarrow T_A = T_E \\ \frac{(1-\alpha_p)S_0}{4} + \sigma T_A^4 = \sigma T_s^4 + F \end{array} \right.$

surface

From the second equation:

$$F = \frac{(1-\alpha_p)S_0}{4} + \sigma T_A^4 - \sigma T_s^4$$

using $T_A = T_E$ and $T_S = T_E + \Gamma z_E$ this becomes

$$F = \frac{(1-\alpha_p) S_0}{4} + \sigma T_E^4 - \sigma (T_E + \Gamma z_E)^4$$

$\uparrow$
 σT_E^4

$$= 2\sigma T_E^4 - \sigma T_E^4 \left(1 + \frac{\Gamma z_E}{T_E}\right)^4 \approx \sigma T_E^4 \left(1 - \frac{4\Gamma z_E}{T_E}\right)$$

$\ll 1$ $\sim 1/2$

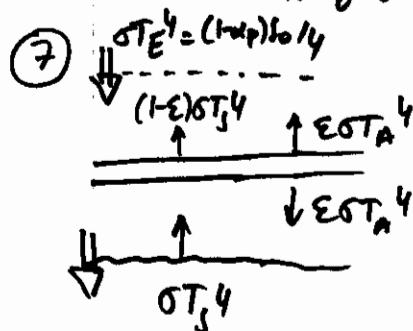
$$F \approx \frac{1}{2} \sigma T_E^4 \sim 120 \text{ W m}^{-2}$$

When $S_0 \uparrow$ by 10%, $T_E \uparrow$ by $(1.1)^{1/4} - 1 = 2\%$

thus $F \uparrow$ by $\{(1.02)^{1/4} \times 120 - 120\} \approx 10 \text{ W m}^{-2}$

$$\textcircled{6} \quad T_{\infty} = \int_0^{\infty} k_A M_A p dz = \underbrace{k_A M_A}_{\text{uniform absorber}} \underbrace{\int_0^{\infty} p dz}_{\text{mass of atmosphere per unit area}}$$

$$\therefore k_A = \frac{T_{\infty}}{M_A \int_0^{\infty} p dz} \approx \frac{2}{6^{-2} \times 10^4} = 2 \cdot 10^{-2} \text{ m}^2 \text{ kg}^{-1}$$



$$(i) \begin{cases} \sigma T_E^4 = (1-\epsilon)\sigma T_S^4 + \epsilon\sigma T_A^4 & (1) \\ \sigma T_E^4 + \epsilon\sigma T_A^4 = \sigma T_S^4 & (2) \end{cases}$$

$$(ii) G = \sigma T_S^4 - [\sigma T_S^4(1-\epsilon) + \epsilon\sigma T_A^4]$$

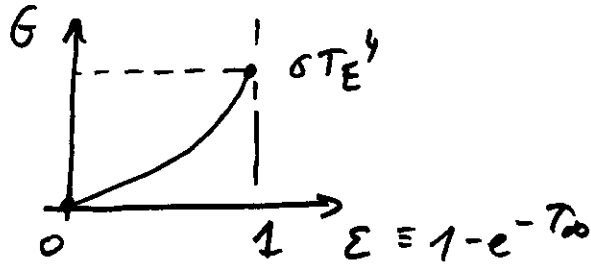
$$= \epsilon\sigma(T_S^4 - T_A^4)$$

Combining (1) & (2) we easily obtain

$$\sigma T_S^4 = \frac{2}{2-\epsilon} \sigma T_E^4$$

$$\sigma T_A^4 = \frac{1}{2-\epsilon} \sigma T_E^4$$

$$\therefore G = \epsilon \left[\frac{2}{2-\epsilon} \sigma T_E^4 - \frac{1}{2-\epsilon} \sigma T_E^4 \right] = \frac{\epsilon}{2-\epsilon} \sigma T_E^4$$



↑
optically
"thin"
= Moon
model

(no greenhouse
effect)

$$T_{\infty} = 0$$

↑
optically
"thick"
= blanket
model

(strongest
greenhouse
effect)

$$T_{\infty} \gg 1$$